

IGSPICE SIMULATION OF INDUCTION MACHINES WITH SATURABLE INDUCTANCES

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Abstract. A new approach for dynamic analysis of induction machines with saturable inductances is presented. The machine dynamics are represented by a set of nonlinear time-varying differential equations. The machine parameters (i.e. saturable inductances) are modeled by closed form nonlinear functions. The equations which define the machine operation are modeled by d and q equivalent circuits. The equivalent circuits are simulated using IGSPICE software.

Key Words: IGSPICE, Induction Machines, Saturable Inductances.

Introduction

Power systems containing induction machines need to be studied in order to determine the interactions of power systems, drive systems and induction machines [1] through [11]. The problems normally studied are as follow: Will the motor start? If so, what is the magnitude of starting current? What is the duration of starting period? Will the motor start without exceeding the motor's thermal rating? What are the impacts of motor's high starting current on the power system load voltages? What is the magnitude of pulsating torque? How many times the motor can be started without degradation to windings' insulation due to high inrush currents? To study these problems, we need an accurate machine dynamic model for the entire operating range and then use the model to evaluate the transient response of starting current, torque, and motor speed.

The inductances which characterize the machine are normally considered constant in conventional models [3], [4], [5]. However, the inductances can vary widely depending on the state of the flux in different parts of the machine. The flux levels are in turn determined by the machine currents which depends on the operating modes [1]-[11]. This requires that the machine be represented by differential equations with time-varying and nonlinear parameters. The nonlinear time-varying parameters will characterize the saturated and unsaturated modes of the machine operation.

The unsaturated machine models have been routinely solved either on an analog or a digital computer [3], [4], [5]. Methods have been described in the literature for modeling saturation of air gap or magnetizing inductances [6]-[11]. In a recent pa-

per, Lipo and Consoli [1], proposed a method for modeling the stator and rotor leakage reactance saturation. They experimentally measured the stator current and rotor speed of a machine during free acceleration from test and then simulated the machine model with their proposed saturated leakage reactances using an analog computer.

This paper presents a new analytical method for developing induction machine models where machine parameters (i.e leakage and magnetizing inductances) may be saturated depending on the operating mode. The degree of inductance saturation is expressed by some nonlinear functions in terms of exciting currents. Furthermore, from saturated machine equations, equivalent circuit models are developed based on the approach proposed by Krause [3] and these circuits are simulated using IGSPICE software. The machine tested and simulated by Lipo and Consoli [1] is studied and the results are compared.

Induction Machine Equations Including Saturation Effects

The transient behavior of an induction machine is generally represented in an orthogonal coordinate system which is either rotating or stationary [3]. The machine performance can be expressed by the following equations, if we assume the angular speed ω of the rotating axes is not specified [3], [4]:

$$v_{qs} = r_s i_{qs} + p \lambda_{qs} + \omega \lambda_{ds} \quad (1)$$

$$v_{ds} = r_s i_{ds} + p \lambda_{ds} - \omega \lambda_{qs} \quad (2)$$

$$v_{os} = r_s i_{os} + p \lambda_{os} \quad (3)$$

$$v_{qr} = r_r i_{qr} + p \lambda_{qr} + (\omega - \omega_r) \lambda_{dr} \quad (4)$$

$$v_{dr} = r_r i_{dr} + p \lambda_{dr} - (\omega - \omega_r) \lambda_{qr} \quad (5)$$

$$v_{or} = r_r i_{or} + p \lambda_{or} \quad (6)$$

where λ is the total flux-linkage of a particular winding, and p is the operator $\frac{d}{dt}$.

The voltages v_{qs} , v_{ds} , v_{os} , v_{qr} , v_{dr} , and v_{or} are the applied stator and rotor voltages. The equations of transformation relating the $d-q-o$ stator and rotor voltages to the actual applied abc voltages are given in [3]. Since only 3-wire balanced voltage systems without a neutral return are considered, the zero quantities i_{or} , i_{os} , v_{or} , and v_{os} are nonexistent. Therefore the Eq. (3) and Eq. (6) will not be included in this analysis. Furthermore, we will concentrate our analysis on single excited machines, consequently, the machine rotor windings are short and the rotor voltages, v_{qr} , and v_{dr} are identically zero. Assuming that the machine is fed from a balanced voltage set and the reference frame is fixed in the stator [3], then the applied stator

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voltages can be expressed as:

$$v_{qs} = \left(\frac{2}{3}\right)^{1/2} V_{L-L} \cos \omega_e t \quad (7)$$

$$v_{ds} = -\left(\frac{2}{3}\right)^{1/2} V_{L-L} \sin \omega_e t \quad (8)$$

where V_{L-L} is the rated line-to-line rms source voltage and ω_e is the source electrical angular velocity.

The flux linkages λ_{qs} , λ_{ds} , λ_{qr} and λ_{dr} of Eq. (1) through Eq. (4) are related to the currents by

$$\lambda_{qs} = (L_{lsa} + L_{lsi})i_{qs} + L_m(i_{qs} + i_{qr}) \quad (9)$$

$$\lambda_{ds} = (L_{lsa} + L_{lsi})i_{ds} + L_m(i_{ds} + i_{dr}) \quad (10)$$

$$\lambda_{qr} = (L_{lra} + L_{lri})i_{qr} + L_m(i_{qs} + i_{qr}) \quad (11)$$

$$\lambda_{dr} = (L_{lra} + L_{lri})i_{dr} + L_m(i_{ds} + i_{dr}) \quad (12)$$

In the above equations, we have used the modeling procedure proposed by Lipo [1]. In this approach, the total stator and rotor leakage inductances are separated into air-dependent and iron-dependent portions. The terms L_{lsi} and L_{lri} correspond to the sum of the iron-dependent saturated leakage inductances which represent the leakage flux of slot, Zig-Zag, belt and skew leakage for stator and rotor respectively. These different components of leakage flux are discussed in [1] - [11]. The terms L_{lsa} and L_{lra} correspond to the air-dependent end-winding leakage inductances and they are assumed to be constant [1]. The term L_m is the magnetizing inductance and is also assumed to be saturable. When the saturation occurs in the machine, the inductances, L_m , L_{lsi} and L_{lri} become nonlinear and the value of these inductances are determined from their exciting currents. Using Eq. (9) through Eq. (12) in Eq. (1) through Eq. (5), then we will have

$$v_{qs} = r_s i_{qs} + L_{lsa} p i_{qs} + p \lambda_{qs}^* + \omega \lambda_{ds}^* + \omega L_{lsa} i_{ds} \quad (13)$$

$$v_{ds} = r_s i_{ds} + L_{lsa} p i_{ds} + p \lambda_{ds}^* - \omega \lambda_{qs}^* - \omega L_{lsa} i_{qs} \quad (14)$$

$$v_{qr} = r_r i_{qr} + L_{lra} p i_{qr} + p \lambda_{qr}^* + (\omega - \omega_r) \lambda_{dr}^* - (\omega - \omega_r) L_{lra} i_{dr} \quad (15)$$

$$v_{dr} = r_r i_{dr} + L_{lra} p i_{dr} + p \lambda_{dr}^* - (\omega - \omega_r) \lambda_{qr}^* - (\omega - \omega_r) L_{lra} i_{qr} \quad (16)$$

where " * " represents the flux leakage nonlinear effects which can be expressed as

$$\lambda_{qs}^* = L_{lsi}^* i_{qs} + L_m^*(i_{qs} + i_{qr}) \quad (17)$$

$$\lambda_{ds}^* = L_{lsi}^* i_{ds} + L_m^*(i_{ds} + i_{dr}) \quad (18)$$

$$\lambda_{qr}^* = L_{lri}^* i_{qr} + L_m^*(i_{qs} + i_{qr}) \quad (19)$$

$$\lambda_{dr}^* = L_{lri}^* i_{dr} + L_m^*(i_{ds} + i_{dr}) \quad (20)$$

In Eq. (17) through Eq. (20) L_{lsi}^* , L_{lri}^* and L_m^* are nonlinear inductances and they are functions of their respective exciting currents. To bring this nonlinear behavior into focus, we will define the following terms for q - axis

$$\lambda_{mq}^* = L_{mq}^* i_{mq} \quad (21)$$

$$\lambda_{qsi}^* = L_{qsi}^* i_{qs} \quad (22)$$

$$\lambda_{qs}^* = \lambda_{qsi}^* + \lambda_{mq}^* \quad (23)$$

where $i_{mq} = i_{qs} + i_{qr}$.

Since L_{qsi}^* and L_{mq}^* are time-varying inductances and they are functions of their exciting currents i_{qs} and i_{mq} , we need to use the chain rule of derivative for calculating $p \lambda_{qs}^*$ term of equation (13). Applying the chain rule to equation (23), we have

$$\frac{d}{dt}(\lambda_{qs}^*) = \frac{d(\lambda_{qsi}^*)}{d(i_{qs})} \cdot \frac{d(i_{qs})}{dt} + \frac{d(\lambda_{mq}^*)}{d(i_{mq})} \cdot \frac{d(i_{mq})}{dt} \quad (24)$$

let

$$\bar{L}_{qsi} = \frac{d(\lambda_{qsi}^*)}{d(i_{qs})} \quad (25)$$

$$\bar{L}_{mq} = \frac{d(\lambda_{mq}^*)}{d(i_{mq})} \quad (26)$$

Then, Eq. (24) can be written as

$$p \lambda_{qs}^* = \bar{L}_{qsi} \cdot p i_{qs} + \bar{L}_{mq} p i_{mq} \quad (27)$$

The terms \bar{L}_{qsi} and \bar{L}_{mq} are called the incremental inductances and they are the derivative of the hysteresis loop at a given exciting current. The above procedure can also be carried out for the terms λ_{ds}^* , λ_{qr}^* and λ_{dr}^* . Using the above incremental inductance concept, the saturated machine inductances can be calculated and the machine dynamics in saturated and nonsaturated regions can be expressed by the following equations:

$$v_{qs} = r_s i_{qs} + (L_{lqa} + \bar{L}_{lqsi}) p i_{qs} + \bar{L}_{mq} p i_{mq} + v_1 \quad (28)$$

$$v_{ds} = r_s i_{ds} + (L_{ldsa} + \bar{L}_{lds}) p i_{ds} + \bar{L}_{md} p i_{md} - v_2 \quad (29)$$

$$v_{qr} = r_r i_{qr} + (L_{lqra} + \bar{L}_{lqri}) p i_{qr} + \bar{L}_{mq} p i_{mq} + v_3 \quad (30)$$

$$v_{dr} = r_r i_{dr} + (L_{ldr} + \bar{L}_{ldri}) p i_{dr} + \bar{L}_{md} p i_{md} - v_4 \quad (31)$$

and v_1 , v_2 , v_3 , and v_4 can be written as

$$v_1 = (\lambda_{lqsa} + \lambda_{lds}^* + \lambda_{md}^*) \cdot \omega \quad (32)$$

$$v_2 = (\lambda_{lqsa} + \lambda_{lds}^* + \lambda_{mq}^*) \cdot \omega \quad (33)$$

$$v_3 = (\lambda_{ldra} + \lambda_{ldri}^* + \lambda_{md}^*) \cdot (\omega - \omega_r) \quad (34)$$

$$v_4 = (\lambda_{lqra} + \lambda_{lqri}^* + \lambda_{mq}^*) \cdot (\omega - \omega_r) \quad (35)$$

Saturated Flux Models

In order to simulate the induction machine dynamics, the saturable flux leakage and magnetizing flux have to be modeled into some functional forms. The flux saturation can be reasonably modeled by a normal magnetizing curve rather than the hysteresis loop. The functional form used in this study for modeling the flux saturation effect is given by the following equation,

$$\lambda = a_1 \arctan(a_2 i) + a_3 i \quad (36)$$

By using the actual data collected for λ and i , the coefficients a_1 , a_2 , and a_3 can be estimated. the incremental inductance is then

$$L_{inc} = \frac{d\lambda}{di} = \frac{a_1 \cdot a_2}{1 + a_2^2 i^2} + a_3 \quad (37)$$

The techniques to identify the constants in Eq. (36) are based on an nonlinear least squares estimation algorithm developed by Marquardt [12]. this can be done by first writing the Taylor series of λ_{isi}^* through linear terms, i.e.,

$$\lambda_{isi}^*(i_n, X + dX) = \lambda_{isi}^*(i_n, X) + \sum_{j=1}^3 \left(\frac{\partial \lambda_{isi}^*}{\partial x_j} \right) (dx)_j \quad n = 1, \dots, N \quad (38)$$

where

$$X = [a_1, a_2, a_3]^T = [x_1, x_2, x_3]^T$$

The vector X in Eq. (38) represents the constants in Eq. (36) which need to be estimated, and i_n is the exciting current, dX is the small correction vector to X . From Eq. (38), the following matrix equation can be obtained as

$$(F + I\theta)dX = h \quad (39)$$

where

$$F(3 \times 3) = J^T \cdot J$$

$$J(n \times 3) = \left(\frac{\partial \lambda_{isi}^*(i_n)}{\partial x_j} \right), \quad n = 1, \dots, N; \quad j = 1, 2, 3$$

$$h(3 \times 1) = \sum_{i=1}^N [\lambda_{isi}^*(i_i, X + dX) - \lambda_{isi}^*(i_i, X)] \frac{\partial \lambda_{isi}^*(i_i, X)}{\partial x_j}$$

In the above equations, N is the total number of experimental data points. Assuming

$$\Delta \lambda_{isi}^* = \sum_{i=1}^N [\lambda_{isi}^*(i_i) - \hat{\lambda}_{isi}^*(i_i)]^2 \quad (40)$$

where λ_{isi}^* is the measured flux leakage and $\hat{\lambda}_{isi}^*$ is the estimated one, the value of θ in Eq. (39) is chosen based on the following strategy:

Let an arbitrary constant v to be greater than one, and let $\theta^{(r-1)}$ represent the value of θ in Eq. (39) at the iteration step $r - 1$. Initially, let $\theta^{(0)}$ to be a constant, say 10^{-2} . Compute $\Delta \lambda_{isi}^*(\theta^{(r-1)})$ and $\Delta \lambda_{isi}^*(\theta^{(r-1)}/v)$. Then the value of $\theta^{(r)}$

at iteration step r can be selected by the following procedure:

1. If $\Delta \lambda_{isi}^*(\theta^{(r-1)}/v) \leq \Delta \lambda_{isi}^*(\theta^{(r-1)})$, Let $\theta^{(r)} = \theta^{(r-1)}/v$.
2. If $\Delta \lambda_{isi}^*(\theta^{(r-1)}/v) > \Delta \lambda_{isi}^*(\theta^{(r-1)})$, and $\Delta \lambda_{isi}^*(\theta^{(r-1)}) \leq \Delta \lambda_{isi}^*(\theta^{(r-1)})$, Let $\theta^{(r)} = \theta^{(r-1)}$.
3. If $\Delta \lambda_{isi}^*(\theta^{(r-1)}/v) > \Delta \lambda_{isi}^*(\theta^{(r-1)})$, and $\Delta \lambda_{isi}^*(\theta^{(r-1)}) > \Delta \lambda_{isi}^*(\theta^{(r-1)})$, increase $\theta^{(r-1)}$ by successive multiplication of v until for some smallest w , $\Delta \lambda_{isi}^*(\theta^{(r-1)w}) \leq \Delta \lambda_{isi}^*(\theta^{(r-1)})$. Then let $\theta^{(r)} = \theta^{(r-1)w}$.

The above estimating process should be continued by updating the unknown vector X as

$$X^{(r+1)} = X^{(r)} + dX^{(r)}$$

until the error between calculated points and corresponding test points can not be reduced further.

Induction Machine Circuit Models for Use in IGSPIICE

The equations (28) - (35) can be represented by the following d -axis and q -axis equivalent circuits.

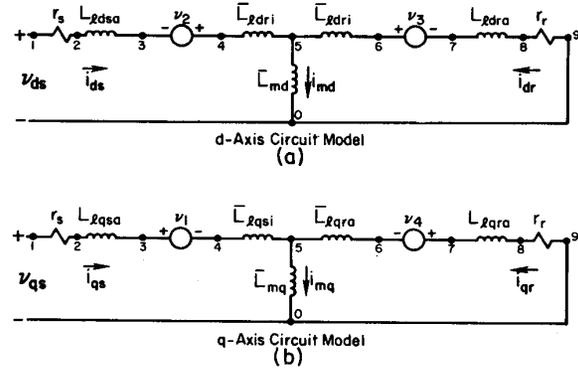


Figure 1: a) The d -axis and b) the q -axis equivalent circuits of a squirrel-cage machine; arbitrary reference frame.

The IGSPIICE is an excellent computer aided analysis package which facilitate fast dynamic simulation of circuit models [13]. To use the IGSPIICE software, one needs to specify the type of a circuit element connected between two nodes. Many types of models exist for simulating active and passive circuit elements. However, the modeling of nonlinear elements are generally left for the user to develop. In the induction machine circuit models, the leakage and magnetizing inductances and the voltage sources v_1, v_2, v_3 and v_4 are current dependent voltage sources which are nonlinear functions of circuit inductances and their exciting currents. The circuit elements (see Fig. 1) connected between nodes (3, 4), (4, 5), (5, 0), (5, 6) and (6, 7) can not be modeled using standard IGSPIICE modules.

To take advantage of the IGSPIICE software, we used fortran subroutines to model the incremental inductances as defined by Eq. (37) and the voltage dependent current sources as defined by equations (3) through (36). Then, these subroutines were linked into the IGSPIICE software for simulation of the electrical transient response of induction machines.

So far, we did not consider the electromagnetic torque equation and mechanical equation relating the total inertia of the machine to the electromagnetic torque. These equations are:

$$T_e = (3/2)(p/2)(\lambda_{ds}^* i_{qs} - \lambda_{qs}^* i_{ds}) \quad (41)$$

$$T_e - T_L = J \left(\frac{d\omega_{rm}}{dt} \right) \quad (42)$$

$$\omega_{rm} = \frac{2}{p} \omega_r \quad (43)$$

where p is the number of the poles, ω_{rm} is the rotor mechanical speed, and J is the rotor inertia.

Equation (42) can be realized by the following circuit topology. The equivalent circuits of Figures 1 and 2 can be used to simulate the electrical and mechanical transient response of induction machines.

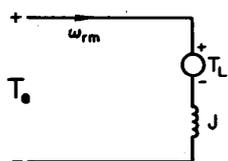


Figure 2: Circuit topology realizing the mechanical equation of the induction machine.

Simulation Results

The machine which was simulated on an analog computer and its response was measured experimentally [1] is studied in this paper. The machine is a three-phase three-wire 230V squirrel-cage machine rated at 5 hp. Table 1 and 2 give the measured line voltages with respect to the exciting currents for the locked-rotor test and the no-load test respectively. The λ_{lrsi}^* and λ_m^* are translated from V_s based on Eq. (44) and Eq. (45), respectively.

Table 1: Locked-Rotor Test Data

| Measured V_s (V, RMS) | Calculated $\lambda_{lrsi}^* = \lambda_{lri}^*$ | Measured I_s (A, RMS) |
|----------------------------|--|----------------------------|
| 0.00 | 0.00000 | 0.00 |
| 6.25 | 0.00677 | 1.88 |
| 15.00 | 0.01625 | 5.00 |
| 24.50 | 0.02653 | 10.00 |
| 32.50 | 0.03520 | 15.00 |
| 40.00 | 0.04332 | 20.00 |
| 47.50 | 0.05144 | 25.00 |
| 54.50 | 0.05902 | 30.00 |
| 60.75 | 0.06579 | 35.00 |
| 67.50 | 0.07310 | 40.00 |
| 72.50 | 0.07852 | 45.00 |
| 77.50 | 0.08393 | 50.00 |
| 82.50 | 0.08935 | 55.00 |
| 87.50 | 0.09476 | 60.00 |
| 91.25 | 0.09882 | 65.00 |
| 95.00 | 0.10289 | 70.00 |

Table 2: No-Load Test Data

| Measured V_s (V, RMS) | Calculated λ_m^* | Measured I_s (A, RMS) |
|----------------------------|-----------------------------|----------------------------|
| 0.00 | 0.00000 | 0.00 |
| 70.00 | 0.15162 | 2.50 |
| 136.25 | 0.29512 | 5.00 |
| 175.00 | 0.37905 | 6.75 |
| 187.50 | 0.40613 | 7.50 |
| 200.00 | 0.43320 | 8.75 |
| 215.00 | 0.46569 | 11.50 |
| 230.00 | 0.49818 | 15.00 |
| 245.00 | 0.53067 | 20.00 |
| 255.00 | 0.55233 | 25.00 |

$$\lambda_{lrsi}^* = \frac{1}{\sqrt{6}} \frac{V_s}{\omega_b} \quad (44)$$

$$\lambda_m^* = \sqrt{\frac{2}{3}} \frac{V_s}{\omega_b} \quad (45)$$

In Eq. (44), λ_{lrsi}^* and λ_{lri}^* are assumed equal, and ω_b is taken to be 377 rad/sec for both Eq. (44) and Eq. (45).

From Table 1 and 2, the coefficients of Eq. (36) which is used to model the flux leakages λ_{lrsi}^* and λ_m^* can be estimated using a computer program which employs the previously discussed algorithm [12], [14]. To do so, a set of initial estimated coefficients must be provided. This can be done by first calculating three different slopes along the saturation curves of λ_{lrsi}^* and λ_m^* . Then the initial estimates can be calculated using the three slopes, namely the increment inductances, and the corresponding exciting currents from Eq. (37).

With the initial estimates and the data provided by Table 1 and 2, the coefficients of Eq. (36) for modelling λ_m^* and λ_{lrsi}^* are estimated. A sample computer output for the coefficients of λ_{lrsi}^* is given by Table 3.

Table 3: Estimated Coefficients of λ_{lrsi}^* .

NONLINEAR LEAST-SQUARES CURVE-FITTING PROGRAM

| | | |
|---------------------------|-------------|-------------------|
| CARD | DEP. VAR. : | MIN Y=0.000E+00 |
| | | MAX Y=1.029E-01 |
| | | RANGE Y=1.029E-01 |
| IND. VAR (I) | NAME | COEF. A (I) |
| 1 | A1 | 2.76848E-02 |
| 2 | A2 | 4.79025E-02 |
| 3 | A3 | 6.74171E-04 |
| NO. OF OBSERVATIONS | | 16 |
| NO. OF COEFFICIENTS | | 3 |
| RESIDUAL ROOT MEAN SQUARE | | 0.00124994 |
| RESIDUAL MEAN SQUARE | | 0.00000156 |
| RESIDUAL SUM OF SQUARES | | 0.00002031 |

Notice in Table 3, the dependent variable is λ_{lrsi}^* and the independent variables are the coefficients of Eq. (36). The RMS error between the estimated λ_{lrsi}^* and the λ_{lrsi}^* listed in Table 1 is given by Table 3 as 0.00125. From Table 3, the flux leakage λ_{lrsi}^* can be written as

$$\begin{aligned} \lambda_{l_{si}}^* &= \lambda_{l_{ri}}^* \\ &= 0.02768 \arctan(0.0479i_{ex}) + 6.74 \times 10^{-4}i_{ex} \end{aligned} \quad (46)$$

Similarly, λ_m^* can be obtained as

$$\lambda_m^* = 0.4095 \arctan(0.1318i_{ex}) \quad (47)$$

Then from Eq. (37) the incremental inductances are given by

$$\bar{L}_{l_{si}} = \bar{L}_{l_{ri}} = \left(\frac{.00133}{1 + .0023i_{ex}^2} \right) + 6.74 \times 10^{-4} \quad (48)$$

$$\bar{L}_m = \left(\frac{0.05396}{1 + 0.0174i_{ex}^2} \right) \quad (49)$$

Note that the functional forms of $\bar{L}_{l_{qsi}}, \bar{L}_{l_{qri}}, \bar{L}_{l_{dsi}}$ and $\bar{L}_{l_{dri}}$ are the same as Eq. (48), except the exciting current i_{ex} will be i_{qs}, i_{qr}, i_{ds} and i_{dr} respectively. Similarly, \bar{L}_{mq} and \bar{L}_{md} terms are the same as Eq. (49) except the exciting current i_{ex} will be equal to i_{mq} and i_{md} respectively. Figures (3), (4), (5), and (6) give the plots of leakage flux, incremental leakage inductance, magnetizing flux, and incremental magnetizing inductance respectively. Figures (7) and (8) show the stator current and the electromagnetic torque when the machine is accelerated at no-load and the saturated parameters and unsaturated parameters are used. The results indicate that the unsaturated model give a substantial error when compared with tests (see Ref. [1]). Therefore, we need to recognize that the model's parameters are time varying and they are a function of exciting current, and in order to obtain the correct simulation results, the saturation effect during the acceleration period can not be neglected.

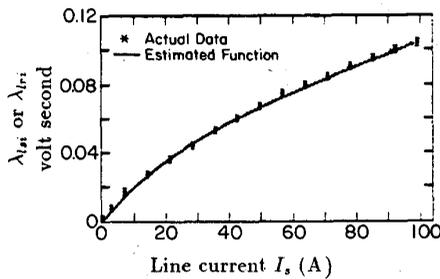


Figure 3: Locked-rotor saturation curve for tested 5-hp machine of reference [1] (I_s is in peak to peak value)

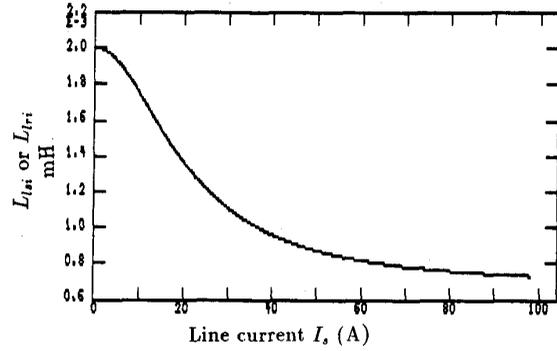


Figure 4: Incremental Inductance $\bar{L}_{l_{si}}$ or $\bar{L}_{l_{ri}}$ (I_s is in peak to peak value)

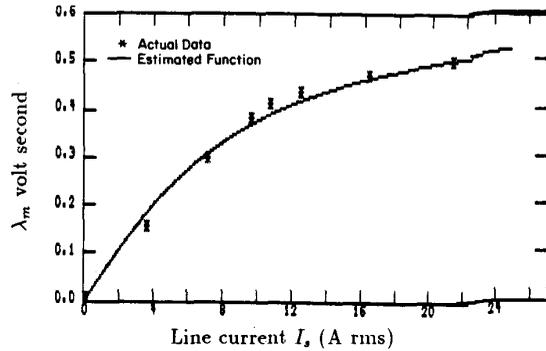


Figure 5: No-load saturation curve for tested 5-hp machine of reference

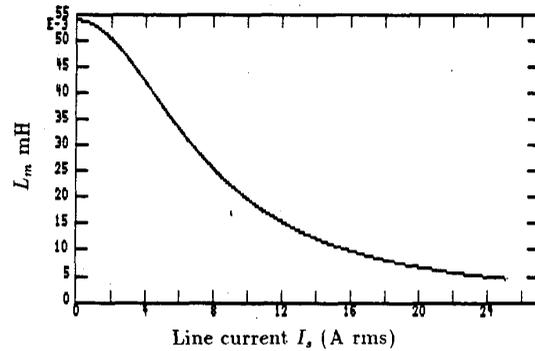


Figure 6: Incremental Inductance \bar{L}_m

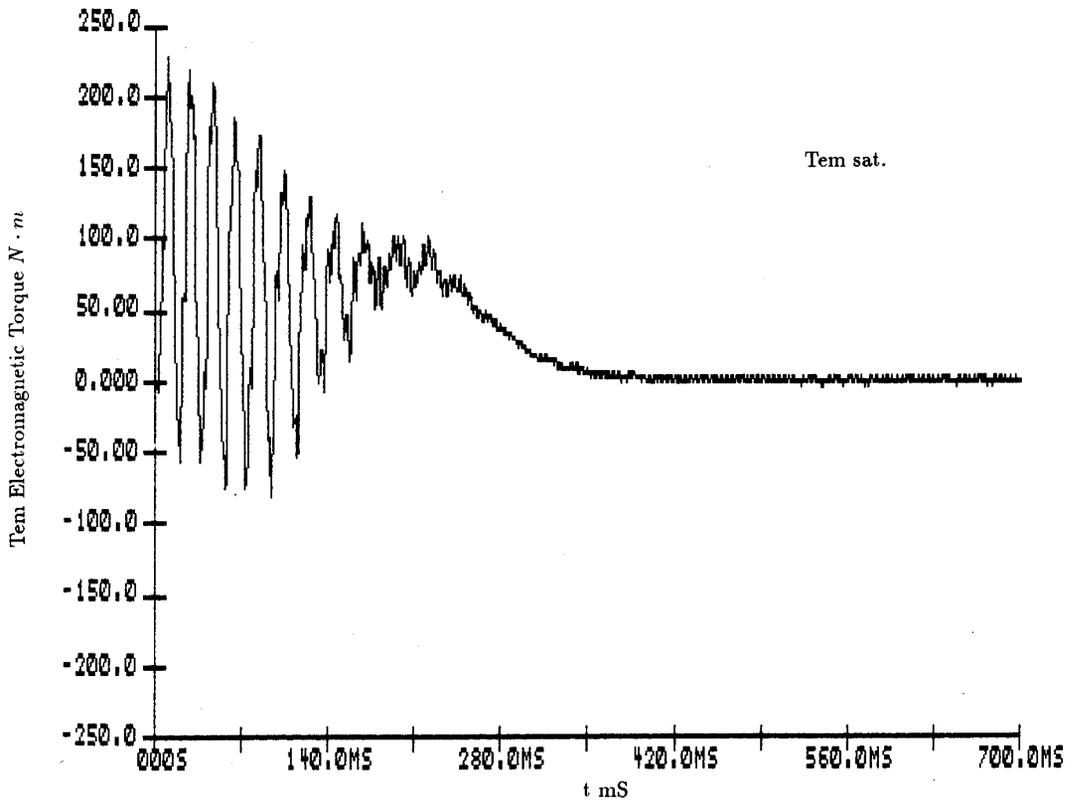
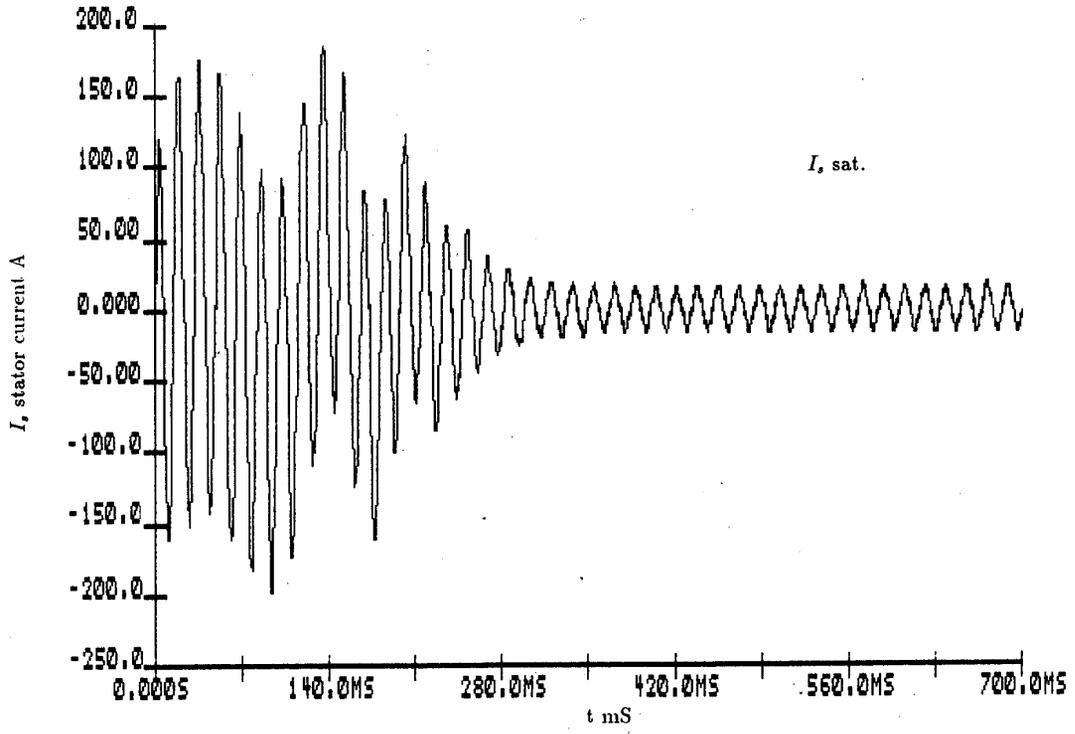


Figure 7: IGSPICE simulation results for free acceleration of test machine using saturated magnetizing and leakage inductances

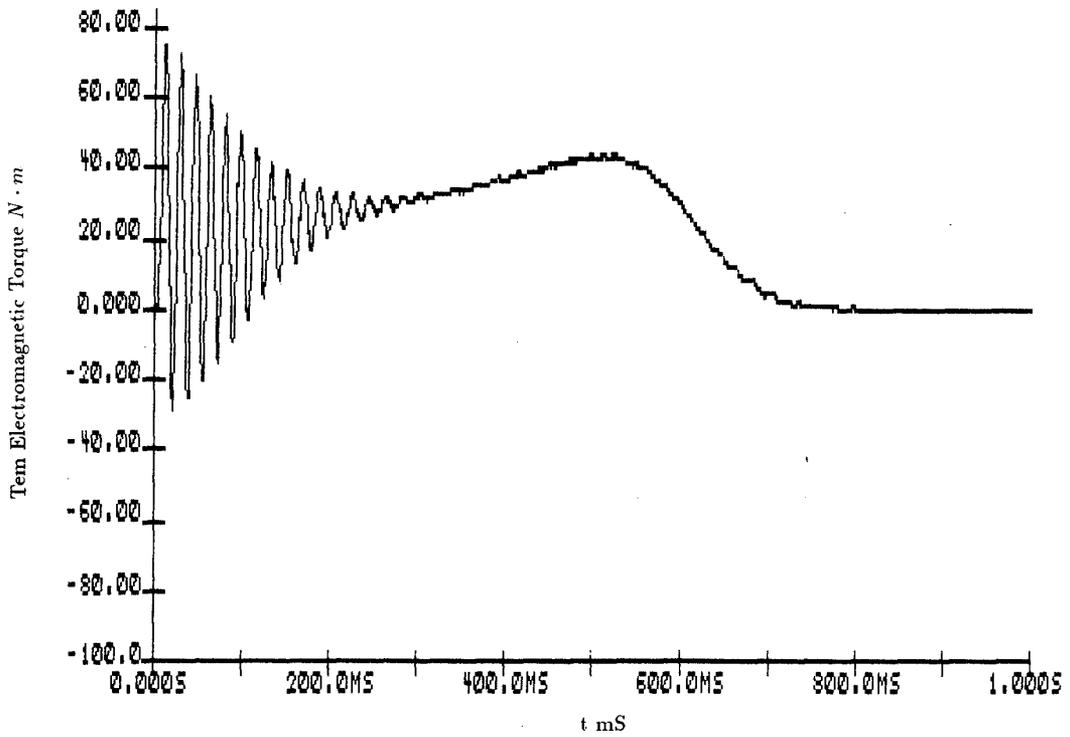
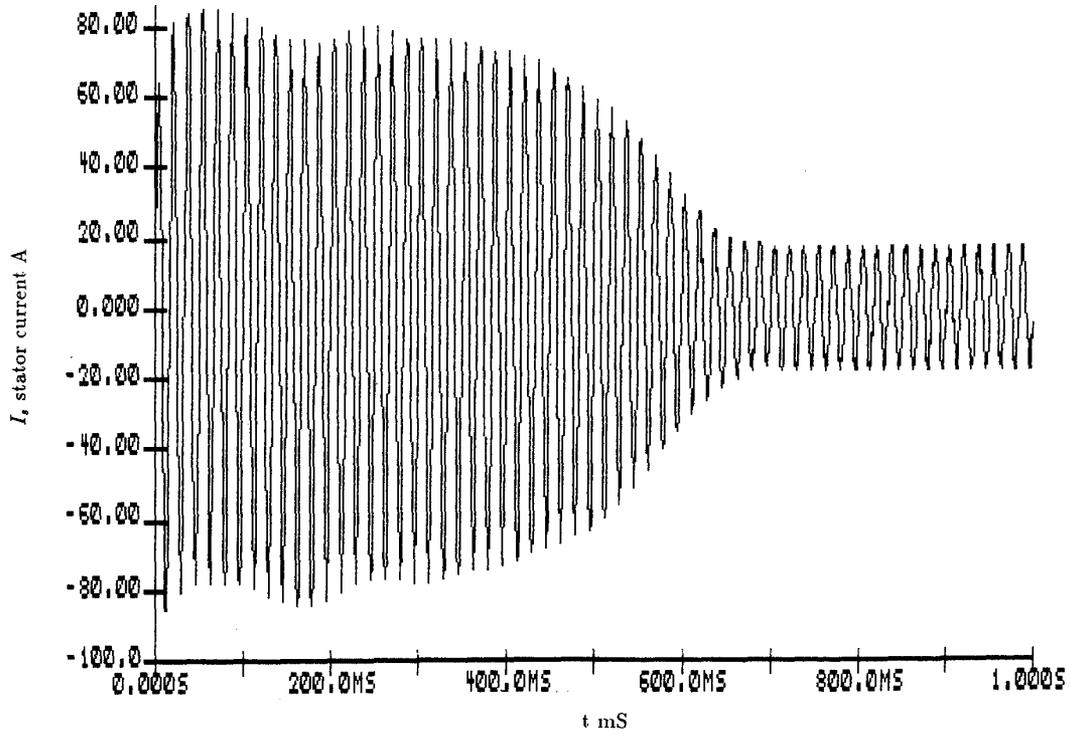


Figure 8: IGSPICE simulation results for free acceleration of test machine using unsaturated (linear) magnetizing and leakage inductances

Conclusion

The paper has presented a new approach for dynamic simulation of induction machines where the machine inductances are nonlinear functions of exciting currents. An induction machine which was tested in reference [1] was modeled and simulated using IGSPICE software. The IGSPICE is a very easy to use and very economical digital computer analysis package. This is in contrast with the analog computer approach which requires many hours of preparation. In addition, since IGSPICE has many circuits modules of active and passive devices, the modeling of drive systems, and power system elements can also be simulated. We believe that IGSPICE will find widespread acceptance with those engineers engaged in the analysis, design and control of industrial drive systems.

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